

The vacuum weighted Einstein field equations: Properties and rigidity of solutions

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Symmetry and Shape

Joint work with Miguel Brozos Vázquez

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CENTRO DE INVESTIGACIÓN
Y TECNOLOGÍA MATEMÁTICA
DE GALICIA

- 1 Context and motivation
- 2 The weighted variational problem
- 3 The vacuum weighted Einstein field equation
 - General properties
 - Isotropic and non-isotropic solutions

Smooth metric measure space

Triple $(M, g, h \, dvol_g)$ where

- M : smooth manifold
- g : semi-Riemannian metric. We focus on Lorentzian metrics.
- $dvol_g$: Riemannian volume element
- $h \in C^\infty(M)$: positive density function ($\nabla h \neq 0$)

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Some notation

- ρ : Usual Ricci tensor ($\rho_{ij} = R^k_{ikj}$)
- Ric: Ricci operator
($\text{Ric}^i_j = g^{ik} \rho_{kj}$)
- τ : Scalar curvature ($\tau = \rho^i_i$)
- Hes_h : Hessian tensor of h
 $\text{Hes}_h(X, Y) = g(\nabla_X \nabla h, Y)$
- Δh : Laplacian of h ($\Delta h = (\text{Hes}_h)^i_i$)

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- Symmetric
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Variational approach:

The Einstein tensor is obtained through a variation of the *Einstein-Hilbert action*:

$$S = \int_{\mathcal{V}} \tau \, d\text{vol}_g$$

The weighted variational problem

Let (M, g) be a Lorentzian manifold and take the action given by the Einstein-Hilbert functional with density:

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- Critical points of this functional under variations

$$g[t] = g + t\delta g, \quad h[t] = h + t\delta h$$

Variations of the metric and its first derivatives vanish at the integration boundary

- We want to preserve the distinguished measure $dV = h \, d\text{vol}_g$, whose variation is $dV[t] = h[t] \, d\text{vol}_{g[t]}$

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$$\delta S_h = \left. \frac{d}{dt} S_h \right|_{t=0} = \int_M \left. \frac{d\tau[t]}{dt} \right|_{t=0} dV + \int_M \tau \left. \frac{d}{dt} \right|_{t=0} dV[t] = 0$$

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The variation of the action reads

$$\delta S_h = \int_M \langle h\rho + \Delta h g - \text{Hes}_h, \delta g \rangle d\text{vol}_g = \int_M \langle D\tau_g^*(h), \delta g \rangle d\text{vol}_g$$

where $\langle T, K \rangle = T^{ij} K_{ij}$, vanishing for all δg at critical points.

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Properties:

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Taking its trace, we have

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Second aim

To understand the geometry of solutions to the vacuum weighted Einstein field equation.

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 - ∇h is **timelike**
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The Jordan form of the **Ricci operator** also plays a role:

$$\text{Type Ia} \\ \text{Ric} = \begin{pmatrix} \alpha_1 & & 0 \\ & \ddots & \\ 0 & & \alpha_n \end{pmatrix}$$

$$\text{Type Ib} \\ \text{Ric} = \begin{pmatrix} a & b & & \\ -b & a & & \\ & & \ddots & \\ & & & \ddots \end{pmatrix}$$

$\{e_1, \dots, e_n\}$ orthonormal basis

$$\text{Type II} \\ \text{Ric} = \begin{pmatrix} \alpha & 0 & & \\ \varepsilon & \alpha & & \\ & & \ddots & \\ & & & \ddots \end{pmatrix}$$

$$\text{Type III} \\ \text{Ric} = \begin{pmatrix} \alpha & 0 & 1 & & \\ 0 & \alpha & 0 & & \\ 0 & 1 & \alpha & & \\ & & & \ddots & \\ & & & & \ddots \end{pmatrix}$$

$\{u, v, e_1, \dots, e_{n-2}\}$ pseudo-ortho. basis
($g(u, u) = g(v, v) = 0, g(u, v) = 1$)

We assume the Jordan form is constant

Lemma

Let $(M, g, h \, d\text{vol}_g)$ be an isotropic solution of the vacuum weighted Einstein field equation. Then Ric is nilpotent and $\Delta h = 0$.

The vacuum weighted Einstein equation reduces to $h\rho = \text{Hes}_h$

Isotropic solutions

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- 1 (M, g) is Ricci-flat and $\text{Hes}_h = 0$
- 2 The Ricci operator is 2-step nilpotent and (M, g) is a Brinkmann wave

Brinkmann wave: (M, g) with a recurrent lightlike geodesic vector field V
($\nabla_X V = \alpha(X)V$, for a 1-form α)

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- 3 The Ricci operator is 3-step nilpotent and (M, g) is a Kundt spacetime

Kundt spacetime: (M, g) with a geodesic lightlike vector field which is

Expansion-free

Shear-free

Twist-free

$$\theta = \frac{1}{n-2} \nabla_i V^i$$

$$\sigma^2 = (\nabla^i V^j) \nabla_{(i} V_{j)} - (n-2)\theta^2$$

$$\omega^2 = (\nabla^i V^j) \nabla_{[i} V_{j]}$$

Theorem

Let (M, g, h) be a locally conformally flat solution.

- 1 If $g(\nabla h, \nabla h) \neq 0$ at a point p , then, on a neighborhood of p , (M, g, h) is locally isometric to a **warped product** $(I \times N, dt^2 \oplus \varphi^2 g^N)$, where
 - N has constant sectional curvature
 - $h(t)$ and $\varphi(t)$ satisfy the following system of ODEs:

$$0 = h'\varphi' - h\varphi'',$$

$$0 = h'' + (n-1)h\frac{\varphi''}{\varphi} + \varepsilon\frac{\tau}{n-1}h.$$

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- 2 If $g(\nabla h, \nabla h) = 0$ on an open subset $\mathfrak{U} \subset M$, then $(\mathfrak{U}, g|_{\mathfrak{U}})$ is a **plane wave** with the metric

$$g(u, v, x_1, \dots, x_{n-2}) = 2dvdu + F(v, x_1, \dots, x_{n-2})dv^2 + \sum_{i=1}^{n-2} dx_i^2,$$

$$\text{where } F(v, x_1, \dots, x_{n-2}) = -\frac{h''(v)}{(n-2)h(v)} \sum_{i=1}^{n-2} x_i^2 + \sum_{i=1}^{n-2} b_i(v)x_i + c(v).$$

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$$\left. \begin{array}{l} g \text{ is positive definite on } \nabla h^\perp \\ \operatorname{Ric} \text{ is self-adjoint} \\ \operatorname{Ric}(\nabla h) = \lambda_1 \nabla h \end{array} \right\} \Rightarrow \operatorname{Ric} \text{ is diagonalizable (Type Ia)}$$

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- If ∇h is spacelike, Ric **does not diagonalize** in general
- We focus on 4-dimensional solutions

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- 2 (M, g) is a **Kundt spacetime** and, depending on the causal character of ∇h , one of the following applies:
 - 1 If $g(\nabla h, \nabla h) = 0$, then Ric is nilpotent and ∇h determines the lightlike parallel line field. Moreover, if Ric vanishes or is 2-step nilpotent, the underlying manifold is a **pp-wave**.

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 - 1 If $g(\nabla h, \nabla h) = 0$, then Ric is nilpotent and ∇h determines the lightlike parallel line field. Moreover, if Ric vanishes or is 2-step nilpotent, the underlying manifold is a *pp-wave*.
 - 2 If $g(\nabla h, \nabla h) \neq 0$, then ∇h is spacelike and the distinguished lightlike vector field is orthogonal to ∇h .

Sketch of the proof (non-isotropic case)

$$\text{Type Ia} \\ \text{Ric} = \begin{pmatrix} \lambda_1 & 0 & 0 & 0 \\ 0 & \lambda_2 & 0 & 0 \\ 0 & 0 & \lambda_3 & 0 \\ 0 & 0 & 0 & \lambda_4 \end{pmatrix}$$

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$\{e_1, \dots, e_n\}$ orthonormal basis with $e_1 = \nabla h / |\nabla h|$

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- **Type Ia:** We analyze the geometric structure according to the **number of distinct eigenvalues** of Ric
 - $\lambda_2, \lambda_3, \lambda_4$ cannot be pairwise distinct \rightsquigarrow Multiply warped product

Sketch of the proof (non-isotropic case)

$$\text{Type Ia} \\ \text{Ric} = \begin{pmatrix} \lambda_1 & 0 & 0 & 0 \\ 0 & \lambda_2 & 0 & 0 \\ 0 & 0 & \lambda_3 & 0 \\ 0 & 0 & 0 & \lambda_4 \end{pmatrix}$$

$$\text{Type Ib} \\ \text{Ric} = \begin{pmatrix} \lambda & 0 & 0 & 0 \\ 0 & a & b & 0 \\ 0 & -b & a & 0 \\ 0 & 0 & 0 & \alpha \end{pmatrix}$$

$\{e_1, \dots, e_n\}$ orthonormal basis with $e_1 = \nabla h / |\nabla h|$

- **Type Ia:** We analyze the geometric structure according to the number of distinct eigenvalues of Ric
 - $\lambda_2, \lambda_3, \lambda_4$ cannot be pairwise distinct \rightsquigarrow Multiply warped product
- **Type Ib:** Use $\text{div } R = 0$ and the Einstein equations to obtain information on the eigenvalues, the curvature and the Christoffel symbols
 - Polynomial system on 5 variables
 - Show $b = 0$
 - There are solutions with $\text{div } R \neq 0$

Non-diagonalizable case with real eigenvalues

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The 2 remaining normal forms for Ric:

$$\text{Type II} \\ \operatorname{Ric} = \begin{pmatrix} \lambda & 0 & 0 & 0 \\ 0 & \alpha & 0 & 0 \\ 0 & \varepsilon & \alpha & 0 \\ 0 & 0 & 0 & \beta \end{pmatrix}$$

$$\text{Type III} \\ \operatorname{Ric} = \begin{pmatrix} \lambda & 0 & 0 & 0 \\ 0 & \alpha & 0 & 1 \\ 0 & 0 & \alpha & 0 \\ 0 & 0 & 1 & \alpha \end{pmatrix}$$

$\{\nabla h, u, v, e_1\}$ pseudo-orthonormal basis ($g(u, u) = g(v, v) = 0, g(u, v) = 1$)

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 - We classified 4D *pr*-wave solutions.

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- **Type III:** solutions are **Kundt spacetimes** with geodesic vector field u .

Thank you!



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